Lecture 3 Ehnesmann connections

In the previous lecture we introduced principal G-bundles : $G \rightarrow P \xrightarrow{\pi} M$ is a fibre bundle whose fibres are G-torsors & whose local trunalisations $\pi^{-1}U \rightarrow U \times G$ are G-equivariant.

Let $P \xrightarrow{=} M$ be a principal G-bindle. Then let $p \in P$ and consider the derivative $(\pi_*)_p: T_p P \longrightarrow T_{mp}M$, which is a surjective linear map. Its bornel $V_p := her(\pi_*)_p$ is called the vertical subsequee. A vector field $\xi \in \mathfrak{E}(P)$ is vertical if $\xi(p) \in V_p$ $\forall p \in P$. The lie brachet of two vertical vector fields is vertical. The vertical subseques span a G-invariant integrable distribution $V \subset P$. In deed, since $\pi \circ r_g = \pi$, $\pi_* \circ (r_g)_* = \pi_*$ and hence $(r_g)_* \vee_p = V_{pg}$.

<u>11. Definition</u> An Ehnesmann connection on P is a smooth choice of horizontal subspaces $H_p \subset T_p P$ complementary to V_p : $T_p P = H_p \oplus V_p$ and such that $(r_g)_* H_p = H_{pg}$. In other words, an Ehnesmann connection is a G-invariant distribution $H \subset TP$ complementary to V.

12. Example A G-invariant <u>newannian</u> metric on E defines an Ethnesmann connection by Hp=Vp. This is a familiar construction in Kaluza- Klein Heavy.

Let $P \xrightarrow{-} M$ be a ppal G-bundle. The Gaction on \mathbb{P} defines a smooth map $q \longrightarrow \mathfrak{X}(\mathbb{P})$ arrigning to every $X \in q$ the fundamental vector field $\xi_X \in \mathfrak{X}(\mathbb{P})$ whose value at p is given by $\xi_X(p) = \frac{d}{dt} \left(p e^{tX} \right) \Big|_{t=0}$ 13. Lenne ξ_X is vertical. Proof $\pi_* \xi_X \Big|_p = \frac{d}{dt} \pi(p e^{tX}) \Big|_{t=0} = \frac{d}{dt} \pi(p) \Big|_{t=0} = 0$.

However, since the Exaction is free, the map $X \mapsto \xi_{\chi}(r)$ is an isomorphism $g \stackrel{\cong}{\to} V_p \quad \forall p \in \mathbb{P}$.

14. Lema
$$(r_{g})_{*} \not \equiv \chi = \not \leq_{Ad_{g},(X)}$$
 for all $g \in G, X \in Q$.
Proof let $p \in \mathbb{P}$. Then,
 $(r_{g})_{*} \not \equiv_{X}(p) = \frac{d}{dt} [r_{g}(pe^{tX})]_{t=0} = \frac{d}{dt} (pe^{tX}g)]_{t=0} = \frac{d}{dt} (pge^{tAd_{g},X})|_{t=0} = \not \leq_{Ad_{g},X}(r_{d})$.
15. Definition The connection one-form of a connection HCTP is the q-valued one-form $\omega \in \Omega^{4}(I;q)$
defined by $\omega(\not \equiv) = \begin{cases} X \quad ij \not \equiv = \not \equiv X \qquad NE: time deep define ω by $(\mathbb{P}(P) - lengendy)$.
16. Proposition The connection one-form ober $r_{g}^{*}\omega = Ad_{g}, \infty$.
Proof let $\not \equiv$ be horizontal. There $(r_{g}^{*}\omega)(\not \equiv) = \omega((r_{g})_{*}\not \equiv)$ but H is G-invariant \mathcal{R} betwee $(r_{g}h \not \equiv b a \ box) = Ad_{g}(X)$
so that $(r_{g}^{*}\omega)(\not \equiv) = \mathcal{N}$ of $(r_{g}^{*})_{*} \not \equiv \Delta d_{g}(X)$.
Note that $(r_{g}^{*}\omega)(\not \equiv) = \omega((r_{g})_{*}\not \equiv)$ as well. Now let $\not \equiv = \not \equiv X = Ad_{g}(X)$
but by Lemma 19. $(r_{g}^{*}\omega)(\not \equiv) = \omega((r_{g})_{*}\not \equiv) \omega = Ad_{g}(X)$.
Conversely, $i \not \omega \in \Omega^{4}(I;q)$ ober $r_{g}^{*}\omega \equiv Ad_{g}(\infty)$ and $\omega(\not \equiv) = X$, then $H = her \omega$ is a convection on \mathbb{P} .
In Physics we are used to gauge fields defining convections on gauge bundles. Gauge fields are local descriptions
 $r_{g}^{*}\omega$ invariantly to the travalization: $(P_{g}^{*}\circ S_{g})(a) = (a, e)$. We may then pull backs ω wire $S_{g}:$
 $A_{g}:= S_{g}^{*}\omega \in \Omega^{4}(U_{d};q)$.$

17. Proposition Let
$$\omega_{\alpha} := Ad_{\alpha} \circ \pi^* A_{\alpha} + g_{\alpha}^* \Theta$$
 where Θ is the left-invariant Maurer-Cartan form on G .
Then $\omega_{\alpha} = \omega_{\pi^{-1} U_{\alpha}}$
Proof we prome this in two steps:

(2) wand wa transform in the same way under the right Gradion :

$$r_{g}^{+}(\omega_{\alpha})_{pg} = Ad_{g_{\alpha}(pg)^{-1}} \cdot r_{g}^{*} \pi^{*} s_{\alpha}^{*} \omega + r_{g}^{*} g_{\alpha}^{*} \theta$$

$$= Ad_{(g_{\alpha}(p)g)^{-1}} \cdot r_{g}^{*} \pi^{*} s_{\alpha}^{*} \omega + g_{\alpha}^{*} R_{g}^{*} \theta \qquad (by Gequivariance of g_{\alpha})$$

$$= Ad_{g^{-1}}g_{\alpha}(p)^{-1} \cdot \pi^{*} s_{\alpha}^{*} \omega + g_{\alpha}^{*} (Ad_{g^{-1}} \cdot \theta) \qquad (since \pi \circ r_{g} \in \pi \quad \& \text{ trans}_{1}^{m} \text{ propertues of } \theta)$$

$$= Ad_{g^{-1}} \cdot (Ad_{g_{\alpha}(p)^{-1}} \circ \pi^{*} s_{\alpha}^{*} \omega + g_{\alpha}^{*} \theta)$$

$$= Ad_{g^{-1}} \cdot (\omega_{\alpha})_{p} \qquad \therefore \quad \omega \notin \omega_{\alpha} \text{ acree everywhere } \circ \pi \pi^{-1}(U_{\alpha}) \qquad \blacksquare$$

But ω is a global one-form on \mathbb{P} , hence $\omega_{\alpha} = \omega_{\beta}$ on $\pi^{-1}(U_{\alpha\beta})$. This ellows us to relate $A_{\alpha} \in A_{\beta}$. Indeed, on $U_{\alpha\beta}$, $A_{\alpha} = S_{\alpha}^{*} \omega_{\alpha} = S_{\alpha}^{*} \omega_{\beta} = S_{\alpha}^{*} \left(Adg_{\beta}(S_{\alpha})^{-1} \circ \pi^{*}A_{\beta} + g_{\beta}^{*}\Theta \right)$ $= Adg_{\alpha\beta} \circ A_{\beta} + g_{\beta\alpha}^{*}\Theta \quad (\text{using that } g_{\beta} \circ S_{\alpha} = g_{\beta} \circ g_{\alpha}^{-1} \circ g_{\alpha} \circ S_{\alpha} = g_{\beta\alpha} \right)$ In the language of matrix lie groups, $g_{\beta\alpha}^{*}\Theta = g_{\beta\alpha}^{-1}dg_{\beta\alpha} = g_{\alpha\beta} dg_{\alpha\beta}^{-1} = -dg_{\alpha\beta}g_{\alpha\beta}^{-1} \quad \text{so that}$ $A_{\alpha} = g_{\alpha\beta}A_{\beta}g_{\alpha\beta}^{-1} - dg_{\alpha\beta}g_{\alpha\beta}^{-1} \quad (g_{\alpha})g_{\alpha} = f_{\alpha} \circ f_{\alpha} \circ f_{\alpha} \circ f_{\alpha}$ Similarly, one can ask how $\{A_{\alpha}\}\$ depends on the choice of local section. If $S'_{\alpha}: U_{\alpha} \to \pi^{-1}(U_{\alpha})$ is another local section, $S'_{\alpha}(\alpha) = S_{\alpha}(\alpha)h_{\alpha}(\alpha)$ \exists smooth $h_{\alpha}: U_{\alpha} \to G$ and then $A'_{\alpha} = Ad_{h_{\alpha}^{-1}} \circ A_{\alpha} + h_{\alpha}^{*} \Theta$ which is again of the form of a gauge transformation (by ha').

$$\frac{1}{2} \left(r_{g} \right)_{*} H_{\overline{a}(p)} = \left(r_{g} \right)_{*} \left(r_{$$

Exercise 2 Let I be a gauge transformation in a pal G-brale P - M. Let 5x denote a findamental vector field for the G-action on P. Show that ξ_x is gauge invariant (ie: $\overline{\Phi}_{\mathbf{x}} \xi_{\mathbf{x}} = \xi_{\mathbf{x}}$) and show that if w is the connection one-form for an Ehresmann connection H, then (€-1)^{\$} w is the connection one-form for H[®].

Let $\{A_a\}$ and $\{A_a^{\clubsuit}\}\$ be the gauge fields corresponding to the Elinesmann connections H and H^{\bigstar} . How are they related? First we need to study the local description of Φ . Since Φ preserves fibres, it makes sense to restrict to $\pi^{-1}U_{\bigstar}$. Applying the trivialisation $\Psi_{\alpha}(\Phi(p)) = (\pi(p), g_{\alpha}(\Phi(p)))$ which defines $\overline{\Phi}_{\alpha} : \pi^{-1}U_{\alpha} \rightarrow G$ by $\overline{\Phi}_{\alpha}(p) = g_{\alpha}(\Phi(p))g_{\alpha}(p)^{-1}$

The first observation is that $\overline{\Phi}_{\alpha}$ is constant on the fibres: $\overline{\Phi}_{\alpha}(Pg) = g_{\alpha}(\overline{\Phi}(Pg)) g_{\alpha}(Pg)^{-1} = g_{\alpha}(\overline{\Phi}(P)g) g_{\alpha}(Pg)^{-1} = g_{\alpha}(\overline{\Phi}(P)) g_{\alpha}(P) = \overline{\Phi}_{\alpha}(P)$ $\overline{\Phi}_{\alpha}(Pg) = g_{\alpha}(\overline{\Phi}(Pg)) g_{\alpha}(Pg)^{-1} = g_{\alpha}(\overline{\Phi}(P)) g_{\alpha}(P) = \overline{\Phi}_{\alpha}(P)$ $\overline{\Phi}_{\alpha}(Pg) = g_{\alpha}(\overline{\Phi}(Pg)) g_{\alpha}(Pg)^{-1} = g_{\alpha}(\overline{\Phi}(P)) g_{\alpha}(P) = \overline{\Phi}_{\alpha}(P)$

and hence it defines a smooth map $\varphi_{\alpha}: U_{\alpha} \rightarrow G$. On overlaps $U_{\alpha\beta} \neq \phi$, we have for all a $\in U_{\alpha\beta}$ and $\rho \in \pi^{-1}(\alpha)$,

$$\begin{aligned} \varphi_{\alpha}(\alpha) &= \Im_{\alpha}(\overline{\Phi}(p)) \Im_{\alpha}(p)^{-1} \\ &= \Im_{\alpha}(\overline{\Phi}(p)) \times \Im_{\beta}(\overline{\Phi}(p))^{-1} \Im_{\beta}(\overline{\Phi}(p)) \times \Im_{\beta}(p)^{-1} \Im_{\beta}(p)^{-1} \\ &= \Im_{\alpha}(\overline{\Phi}(p)) \times \Im_{\beta}(\overline{\Phi}(p))^{-1} \Im_{\beta}(\overline{\Phi}(p)) \times \Im_{\alpha}(p)^{-1} \\ &= \Im_{\alpha}(\overline{\Phi}(p)) \oplus_{\alpha}(p)^{-1} \\ &= \Im_{\alpha}(\overline{\Phi}(p)) \oplus_{\alpha}(p) \oplus_{\alpha}(p)^{-1} \\ &= \Im_{\alpha}(\overline{\Phi}(p)) \oplus_{\alpha}(p) \oplus_{\alpha}(p) \oplus_{\alpha}(p) \\ &= \Im_{\alpha}(\overline{\Phi}(p)) \oplus_{\alpha}(p) \oplus_{\alpha$$

We will see later that $\{\varphi_{\alpha}\}$ defines a section of a fibre bundle AdP on M ano dated to the goal bundle P. modify gos Exercise 3 Show that on U_{α} , $A_{\alpha}^{\Xi} = Ad_{\varphi_{\alpha}} \circ (A_{\alpha} - \varphi_{\alpha}^{*}\Theta) \stackrel{!}{=} \varphi_{\alpha} A_{\alpha} \varphi_{\alpha}^{-1} - d\varphi_{\alpha} \varphi_{\alpha}^{-1}$, which is indeed a gauge transformation.

These gauge transformations are conceptually different than the ones relating Ar and AB on Ups. These ones are "global" objects (sections of Ad I) whereas the ones in one-daps are locally defined { 3 gb} on non-empty vertaps.